

Tunneling in Al/Al₂O₃/Al junctions and its direct link with energy gap and tunneling time across the barrier

Edgar J. Patiño

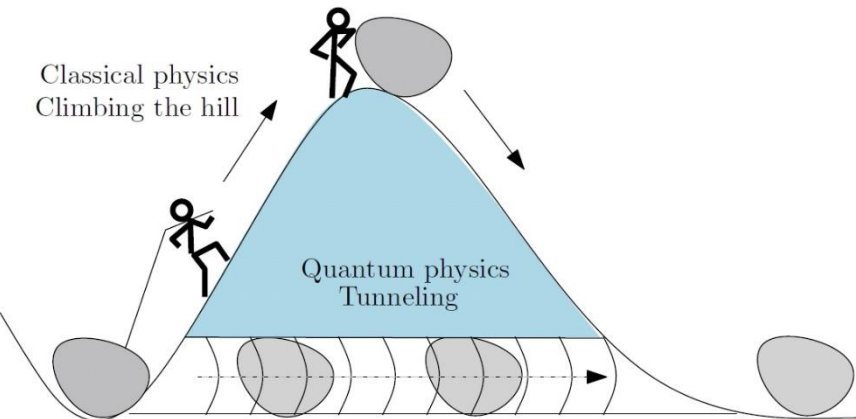
Departamento de Física

Neelima Kelkar

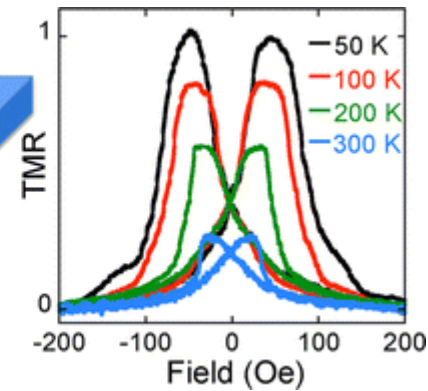
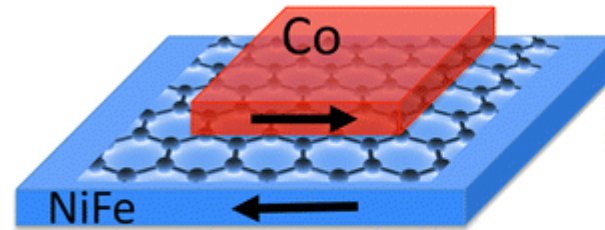
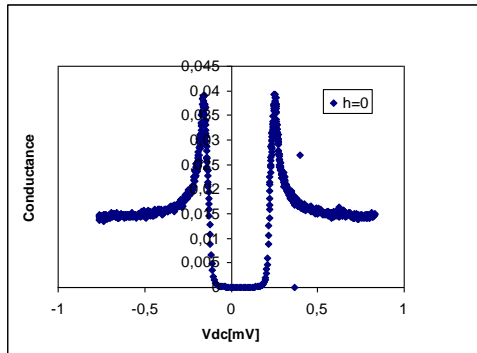
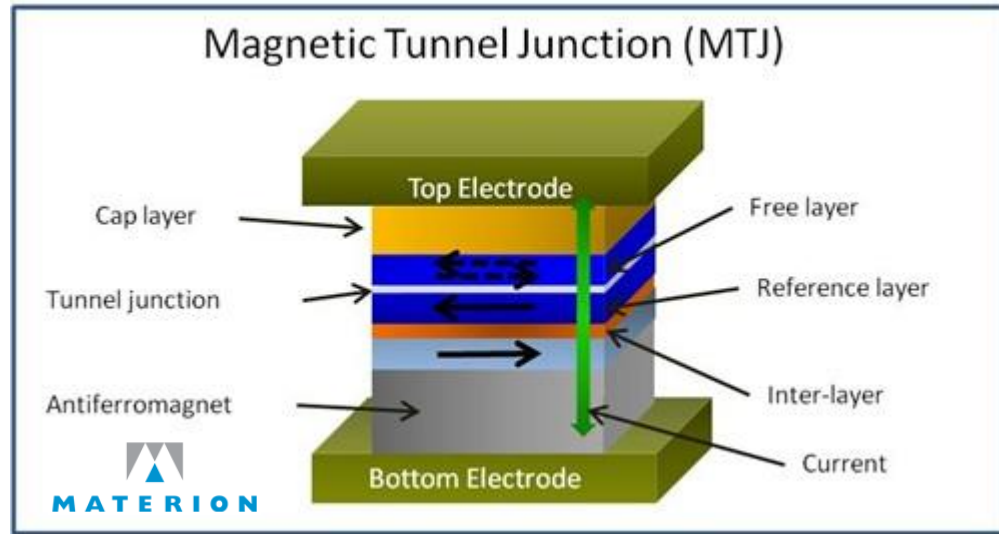


- **Introduction and Motivation**-> Quantum Tunneling.
Why tunneling is important in solid state physics?
Some preliminary works...
- **Experiments Description**-> Additional information
that can be extracted from tunneling experiments
- **Tunneling Dwell time determination**
- **Conclusions**

What is quantum tunneling?

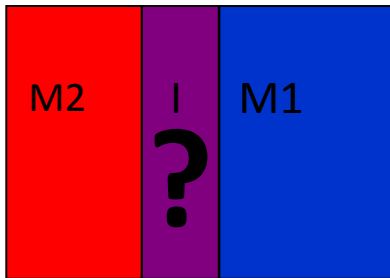
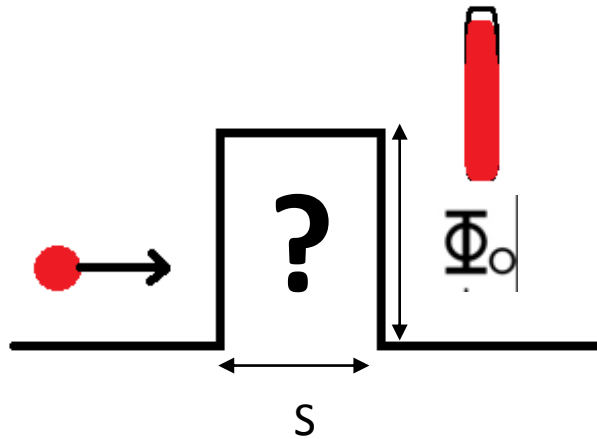


Source: tech for space



Graphene As a Tunnel Barrier: Graphene-Based Magnetic Tunnel Junctions Nano Lett., 2012, 12 (6), pp 3000–3004 (2012)

What people knew before?

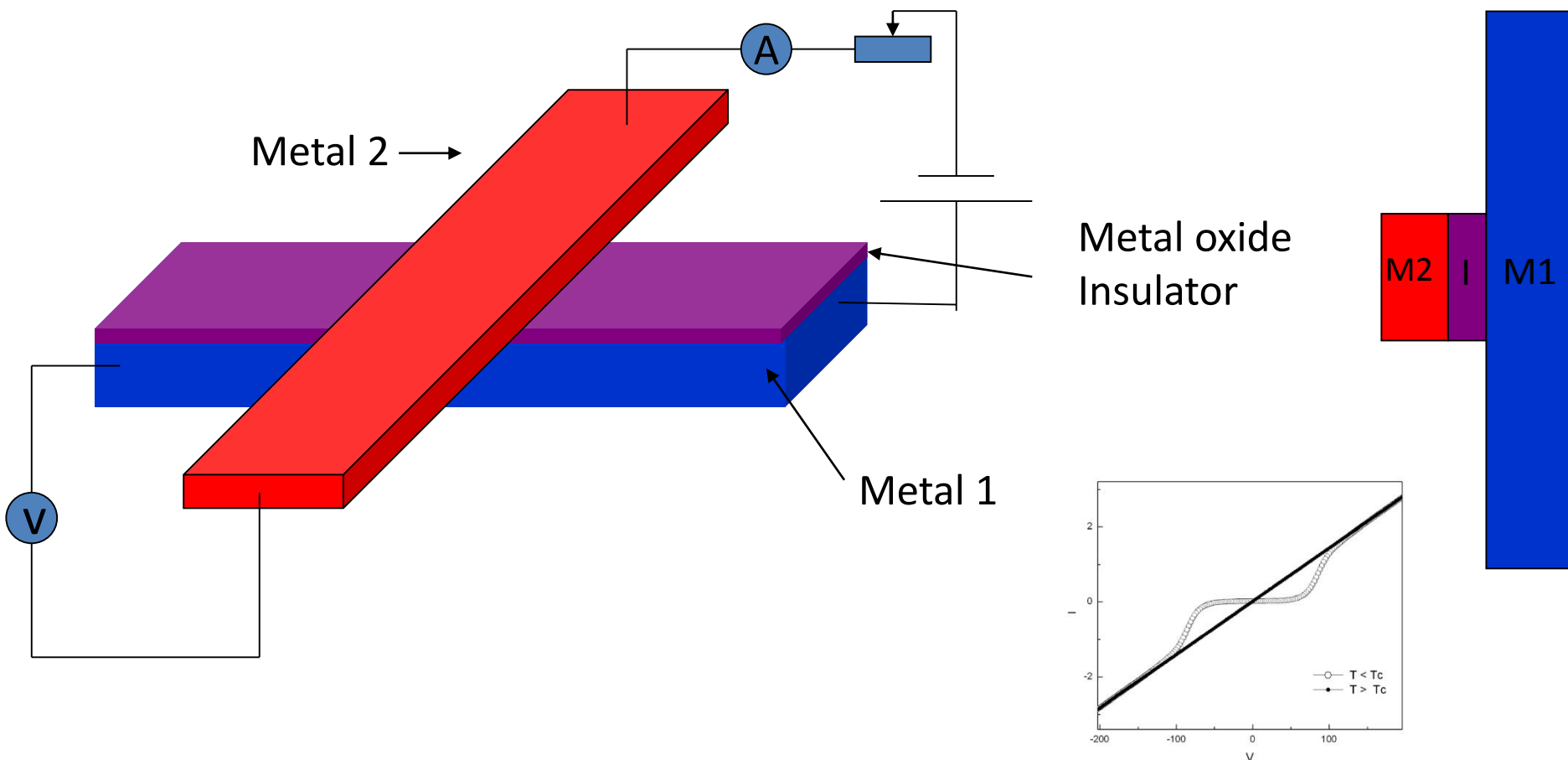


- **Barrier height** decreases with increasing temperature (up to 77 K)....?
- CONTROVERTIAL EXPLANATIONS
 - A) Is Al_2O_3 barrier height temperature dependence?
 - B) Is there a change in the space charge in dielectric?
 - C) Are there trap levels in the insulator?
- Two reports on **barrier width** temperature variation -> NO EXPLANATION (electron effective mass, maybe??)

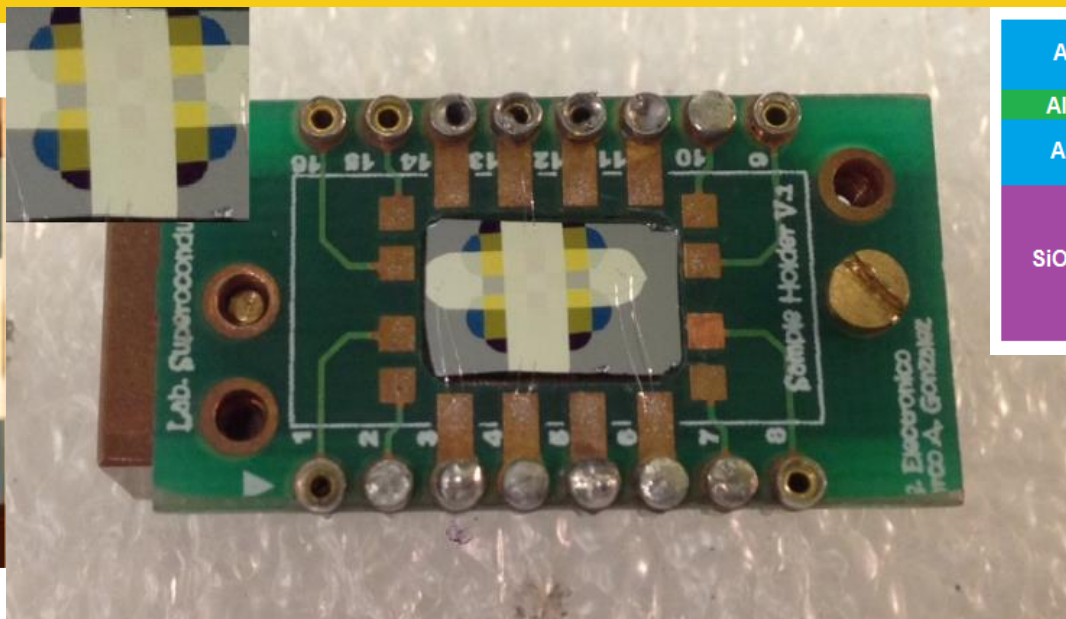
Difficult to produce continuous Al_2O_3 -> pinholes, hot spots and barrier shorts found!

- [1] K. H. Gundlach and A. Wilkinson, Phys. Stat. Sol. (a) **2**, 295 (1970).
- [2] O. L. Nelson and D. E. Anderson, J. Appl. Phys. **37**, 77 (1966)
- [3] J. Kadlec, Solid-State Electronics **17**, 469 (1974).
- [4] V. D. Das and M. S. Jagadeesh, Phys. Stat. Sol. (a) **66**, 327 (1981).
- [5] D. Meyerhofer and S. A. Ochs, J. Appl. Phys. **34**, 2535 (1963).

How to make a solid state tunnel junction ?



These structures can be fabricated utilizing mechanical masks or standard optical lithography techniques



Al (30 nm)
Al ₂ O ₃ (2 nm)
Al (30 nm)
SiO ₂ (~ 300 nm)

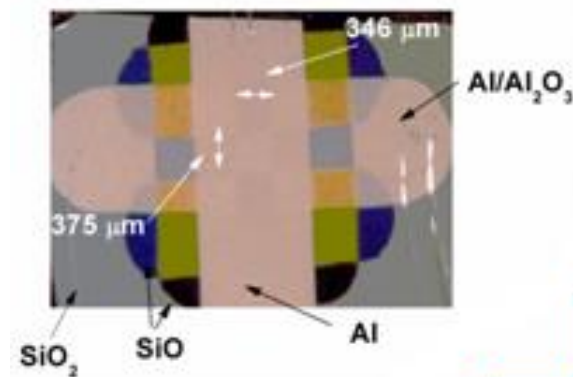
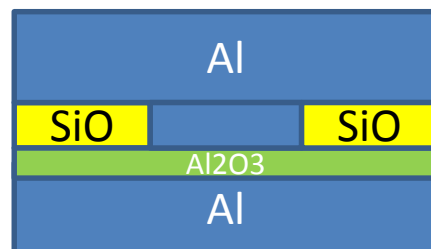
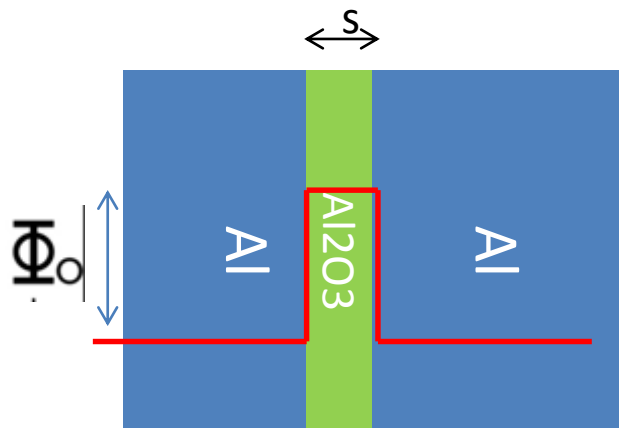
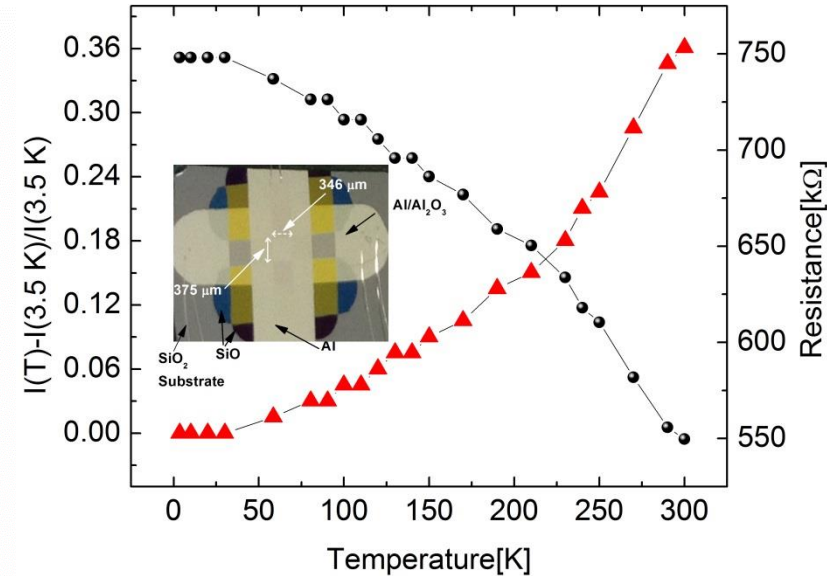
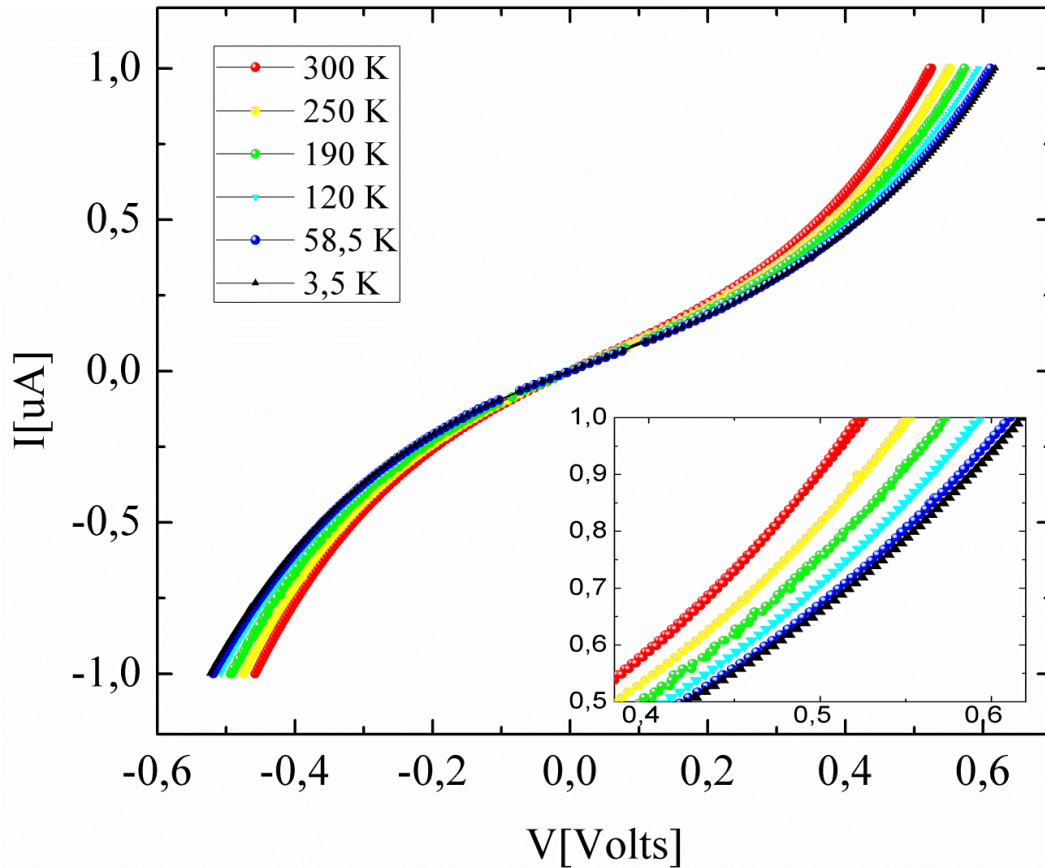


Figure 3: Photograph of planar junctions made using a mask evaporation technique in the system shown in figure. 1.

What are the results?



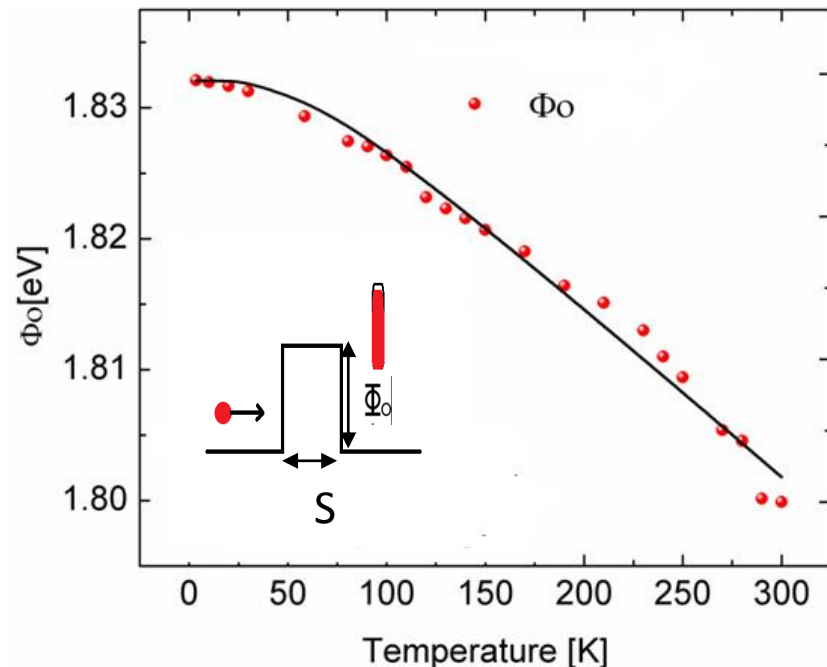
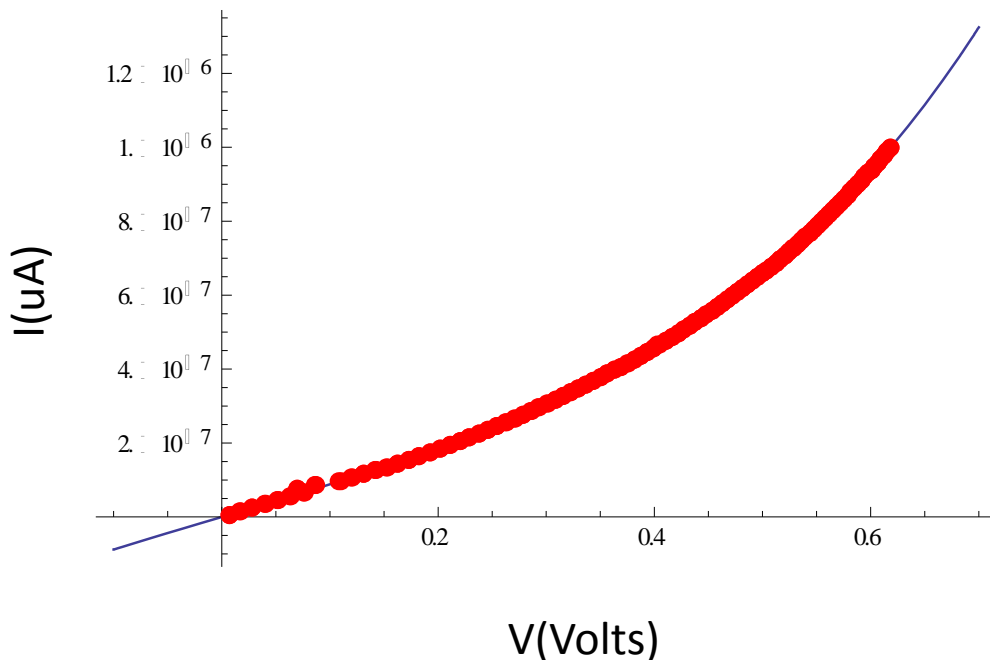
Junction Resistance (triangles) and Tunneling current variation (circles) vs temperature at a bias voltage of 0.5 V

I-V characteristics of Al/Al₂O₃/Al junctions at different temperatures; inset shows zoom in view upper voltages.

Small temperature dependence!

How good is the theoretical fit with the experiment ?

$$J = \left(\frac{e}{2\pi\hbar s^2} \right) \left\{ \left(\varphi_0 - \frac{eV}{2} \right) \exp \left[-\frac{4\pi s}{\hbar} (2m)^{\frac{1}{2}} \left(\varphi_0 - \frac{eV}{2} \right)^{\frac{1}{2}} \right] - \left(\varphi_0 + \frac{eV}{2} \right) \exp \left[-\frac{4\pi s}{\hbar} (2m)^{\frac{1}{2}} \left(\varphi_0 + \frac{eV}{2} \right)^{\frac{1}{2}} \right] \right\}$$



	Estimate	Standard Error
φ_0	2.178	0.00249147
s	19.1074	0.0109113

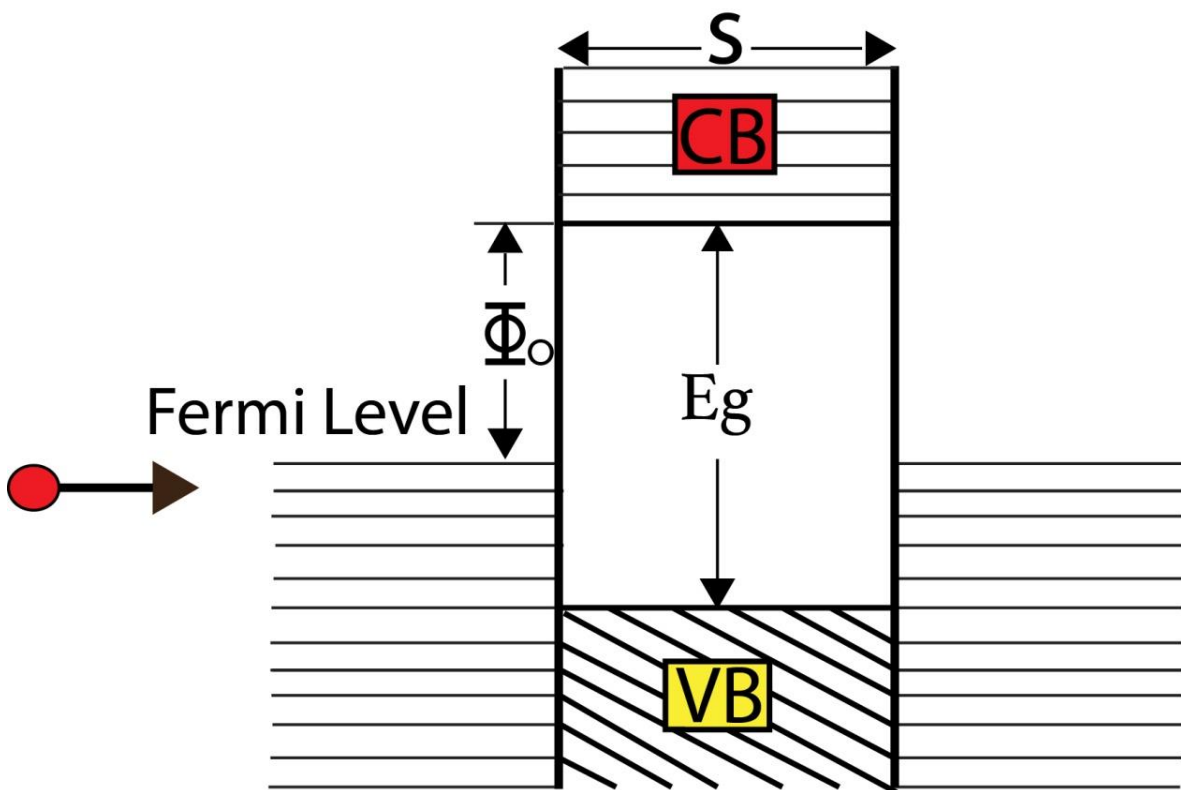
$A \approx 375 \times 346 \text{ (\mu m)}^2$

No pinholes or hot spots!

$S \approx 20.8 \text{ \AA}$

Barrier width \rightarrow Fixed at all Temperatures!

Why is Φ_o changing with temperature ?



Assuming a linear relationship:

$$E_g(T) = \gamma \Phi_o(T)$$

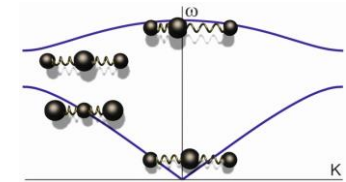
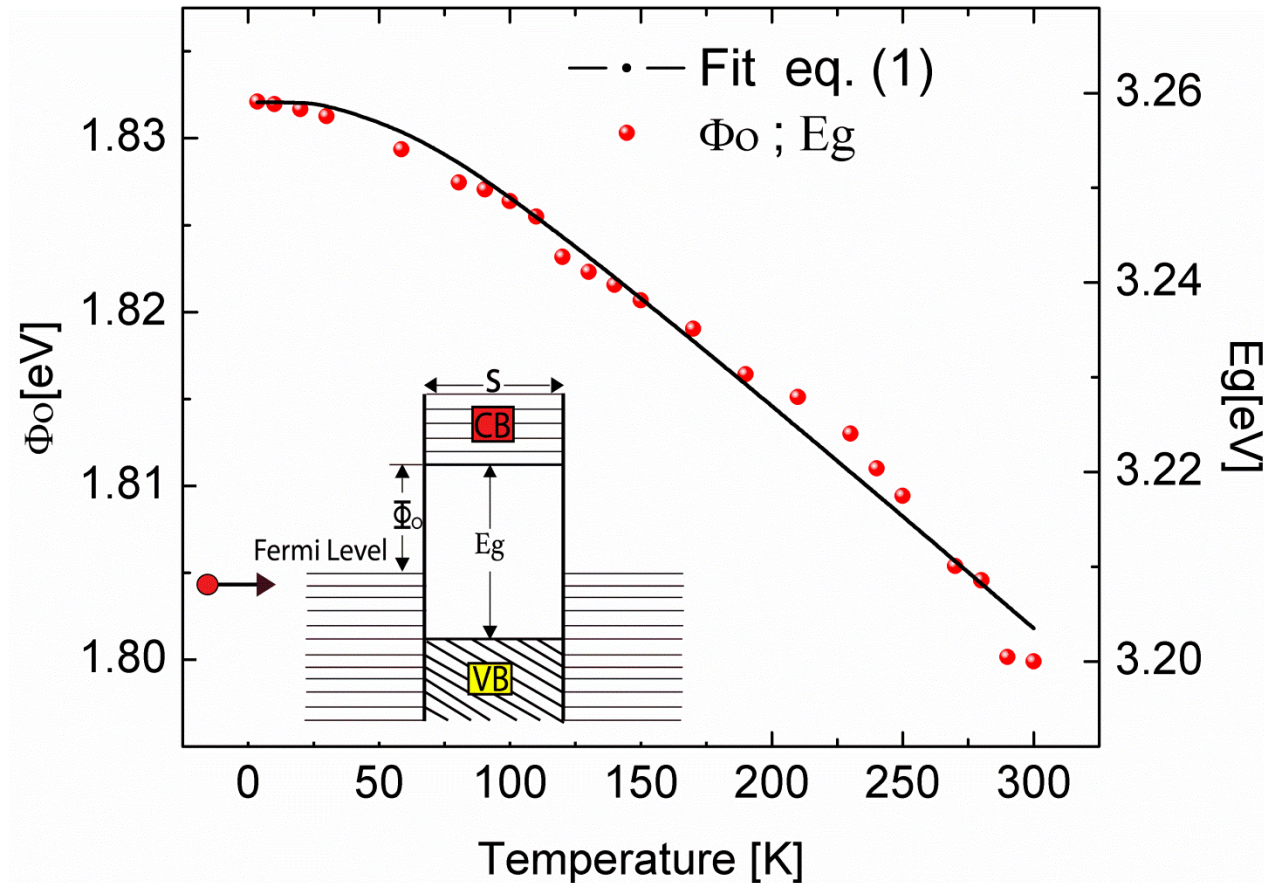
$$\frac{E_g(300 K)}{\Phi_o(300 K)} = \frac{3.2 eV *}{1.8 eV} = 1.7$$

Is E_g changing with temperature as well?

* I. Costina and R. Franchy, Appl. Phys. Lett. 78, 4139 (2001)

Eq. 1

$$E_g(T) = E_g(0) - S \langle \hbar\omega \rangle [\coth(\langle \hbar\omega \rangle / 2kT) - 1]$$



$$E_g(T) = \gamma \phi_0(T)$$

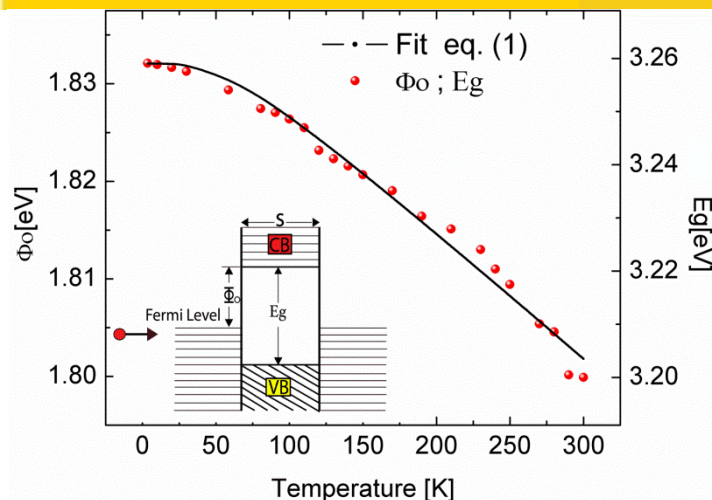
Eq. 1

	Estimate	Standard Error
E _{g0}	3.24348	0.00072269
S	1.16311	0.0613627
$\langle \hbar\omega \rangle$	0.0170526	0.00235379

$$\underline{\text{Avg}[\omega] = 2.05 \times 10^{13} \text{ sec}^{-1}}$$

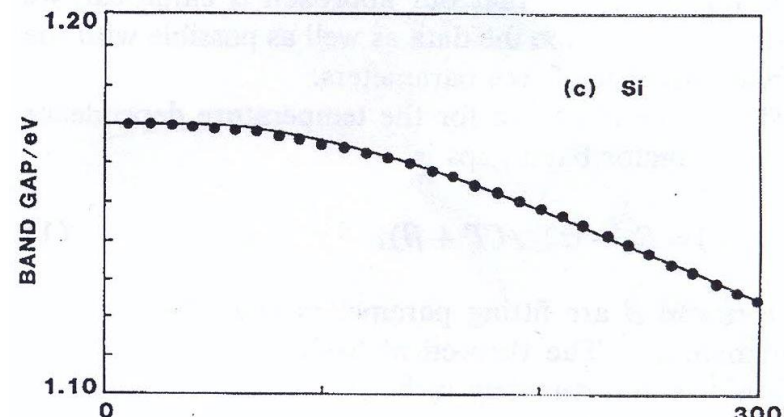
Comparing with speed of sound experiments → $\omega = 2.24 \times 10^{13} \text{ sec}^{-1}$ is obtained !

Comparing with other semiconductors: O'Donnell



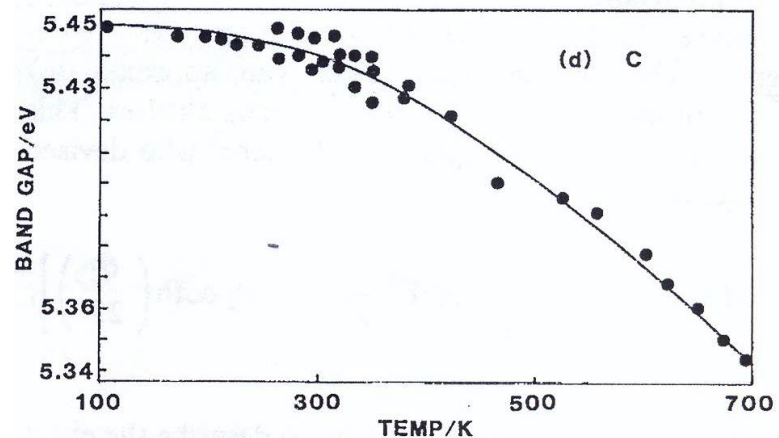
	Estimate	Standard Error
E_{g0}	3.24348	0.00072269
S	1.16311	0.0613627
$\langle \hbar\omega \rangle$	0.0170526	0.00235379

E_g [eV] ↑ Average phonon energy



γ	$E_g(0)$ [eV]	S	$\langle \hbar\omega \rangle$ [meV]	$\langle \omega \rangle$ [$\times 10^{13} \text{sec}^{-1}$]
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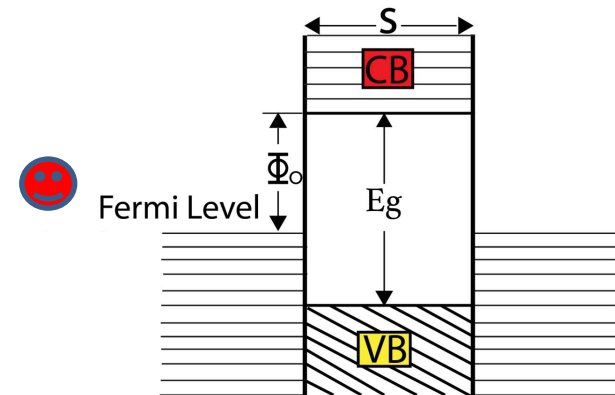
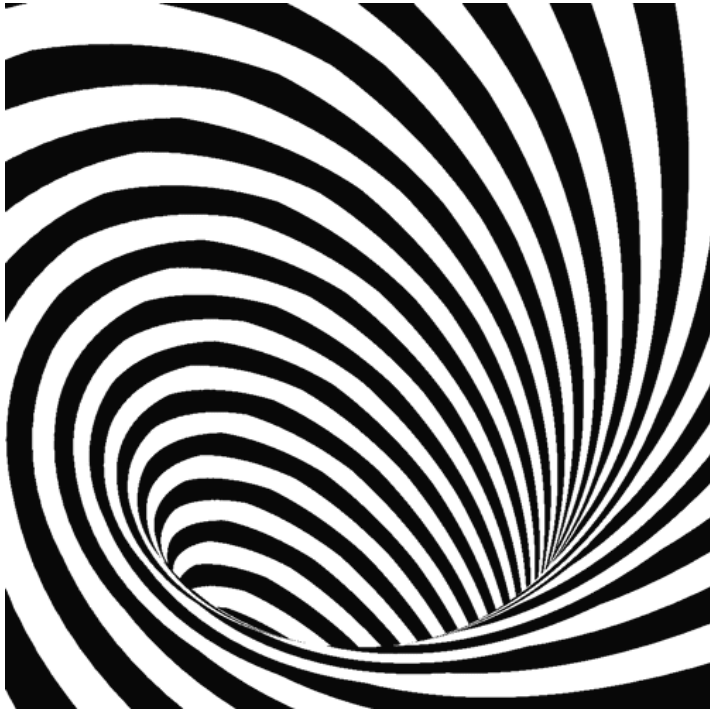
Al ₂ O ₃ Forward bias	1.78	3.26	1.414	13.5	2.05
Al ₂ O ₃ Back bias	1.84	3.24	1.163	17.1	2.59
Si ¹⁴	-	1.17	1.49	25.5	-
GaAs ¹⁴	-	1.52	3.00	26.7	-
GaP ¹⁴	-	2.34	3.35	43.6	-
C ¹⁴	-	5.45	2.31	94.0	-



Tunneling Time

Given that phonon frequencies are
correct we confirm.....

Barrier width “ s ” and height “ ϕ_0 ” are indeed correct!



PHYSICAL REVIEW B, VOLUME 64, 233311

Electron tunneling time measurement by field-emission microscopy

S. K. Sekatskii^{1,2} and V. S. Letokhov^{1,*}

$$\sim 1 \times 10^{-15} \text{ s}$$

LETTER

doi:10.1038/nature11025

Resolving the time when an electron exits a tunnelling barrier

Dror Shafir^{1*}, Hadas Soifer^{1*}, Barry D. Bruner¹, Michal Dagan¹, Yann Mairesse², Serguei Patchkovskii³, Misha Yu. Ivanov^{4,5}, Olga Smirnova⁵ & Nirit Dudovich¹

REPORTS

Attosecond Ionization and Tunneling Delay Time Measurements in Helium

P. Eckle,¹ A. N. Pfeiffer,¹ C. Cirelli,¹ A. Staudte,² R. Dörner,³ H. G. Muller,⁴ M. Büttiker,⁵ U. Keller¹

It is well established that electrons can escape from atoms through tunneling under the influence of strong laser fields, but the timing of the process has been controversial and far too rapid to probe in

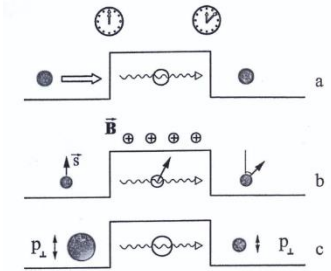


FIG. 1. The idea of the Larmor clock tunneling time measurement. One should invent a clock that starts to measure at the beginning of the barrier crossing and finishes at the end (a). Particle spin vector s rotation in the perpendicular magnetic field B , superimposed in the barrier region (b), or transversal momentum suppression (c) can be used as such a clock.

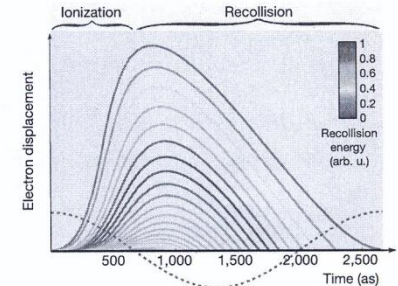


Figure 1 | Electron trajectories contributing to the recollision process. The coloured lines represent the spatio-temporal description of various trajectories; each colour encodes a recolliding energy, increasing from red to blue. The black dashed line shows the electric field along the cycle. arb. u., arbitrary units.

We measured a weighted intensity-averaged tunneling delay time of 6.0 as with a standard deviation of the weighted mean of 5.6 as
 $\sim 1 \times 10^{-18} \text{ s}$

Dwell time (tiempo de habitabilidad)

$$\tau_D(E) = \frac{\int_{x_1}^{x_2} |\Psi(x)|^2 dx}{j}, \quad \text{Eq. 1}$$

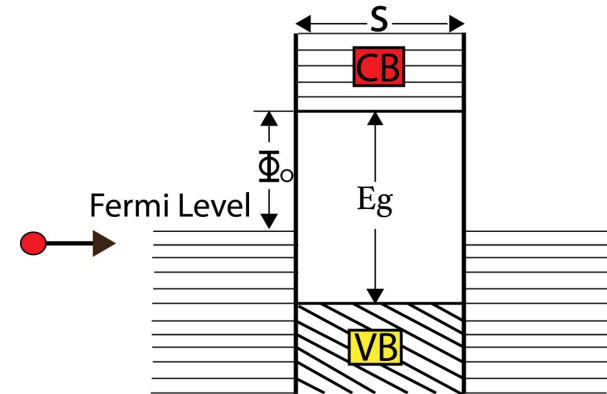
x_1, x_2 , are the classical turning points

momentum



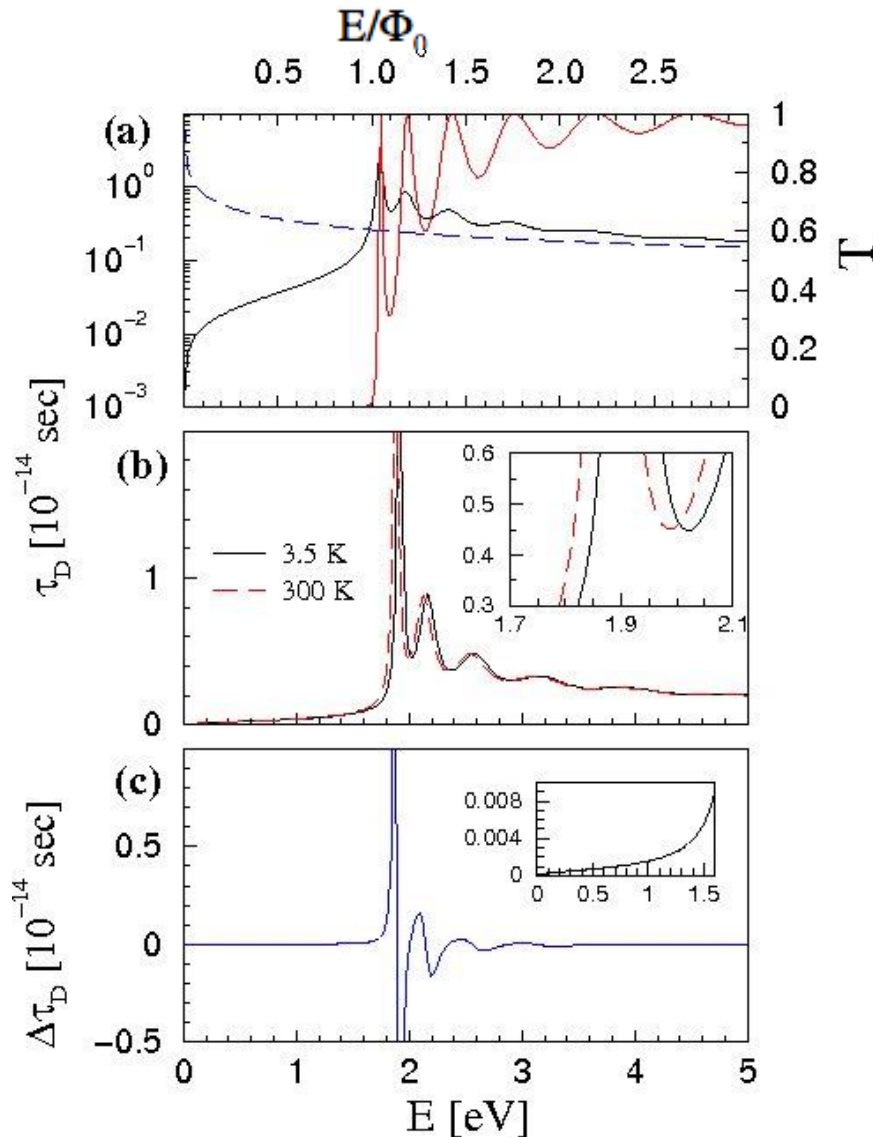
$$j = \hbar k / \mu \quad \leftarrow \text{reduced mass}$$

$$k = \sqrt{2mE} / \hbar$$



Finding the wave function and solving eq. 1 for a rectangular barrier of fixed height and width for $E < \Phi_0$

Finding the wave function and solving eq. 1 for a rectangular barrier :



Average dwell times spent by tunneling electrons within the potential barrier

(a) As a function of the energy divided by the barrier height. The trans-mission coefficient (red line- right scale)

(b) As a function of the energy for two different temperatures.

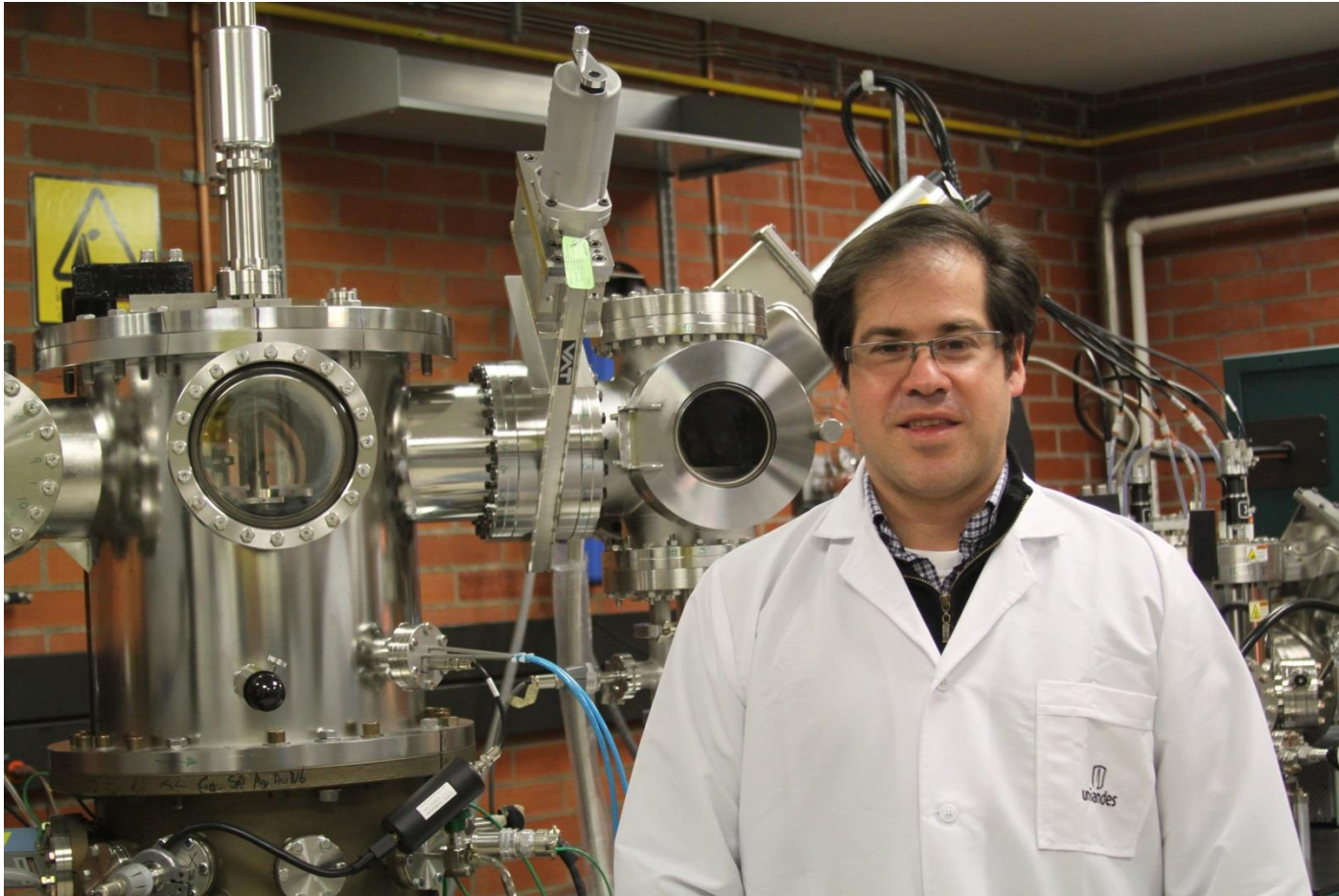
(c) Difference between the dwell time curves at 3.5 and 300 K. Pronounced mostly 300K 3.5K in the resonance regions $\Delta\tau_D = \tau_D(3.5\text{K}) - \tau_D(300\text{K})$

The trans-mission coefficient (red line with scale on right side in (a)).

$$\tau_D = 3.6 \times 10^{-16} \text{ sec at mid barrier energies !}$$

Conclusions

- Tunneling experiments demonstrate a clear temperature dependence of the barrier height
- The barrier height temperature dependence is directly linked to energy gap of the semiconductor BUT barrier width $S \approx 20.8 \text{ \AA}$ does not change.
- The phonon average frequency extracted $\omega = 2.24 \times 10^{13} \text{ sec}^{-1}$ is very close to the one obtained from speed of sound experiments, proving this as an accurate technique .
- Tunneling time determined to be $3.6 \times 10^{-16} \text{ sec}$ at mid-barrier energies
- Tunneling experiments in **other** thin semiconducting materials should provide useful information on energy gap and phonon spectrum



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